



INFLUENCE OF SAMARIUM ION SUBSTITUTION ON STRUCTURAL, ELECTRICAL, AND MAGNETIC PROPERTIES OF CO-MG NANO-FERRITES

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Abstract:

$\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00 \leq x \leq 0.025$) has been prepared using the Sol gel combustion method from powders Of Co, Mg, Fe, and Sm nitrates, which were used as starting materials. To better understand the structure morphology properties, XRD, FESEM, HRTEM, FTIR, Ac and Dc conductivity and magnetic property analyses were done. An x-ray diffraction (XRD) investigation confirmed the creation of a spherical structure with a crystal size of between 33 and 39 nm. It was revealed that the increased doping concentration of Samarium (Sm^{3+}) increases the lattice parameters of the crystal. The X-ray density increases as the concentration of Sm^{3+} doping increased. X-ray density increases due to the Sm^{3+} ion's larger molar mass than the Fe^{3+} ion. Analysis by FESEM indicated that the particles have an irregular size and a high concentration of open pores. It was discovered that ferrite nanoparticles have a spherical in structure with an average size of 75nm can be seen clearly by TEM. According to the FTIR spectra, the primary absorption bands associated with the tetrahedral and octahedral metal-oxygen bonds are located in the region of ($350\text{-}4000\text{cm}^{-1}$). Dielectric spectroscopy revealed Maxwell-Wagner interfacial polarisation, which diminishes with increasing applied field frequency. The sample has a dielectric loss of less than 8, indicating it may be used to fabricate multilayer inductors. The decreases in g-value are well correlated with the variation of particle size. The hopping model explained why the ac conductivity of the samples rose with increasing doping and frequency. Samples with samarium contents of 0.015 and 0.000 showed high impedance stability at room temperature, rendering them suitable for use as a high-frequency. A rise in Sm^{3+} concentration increased M-H loop saturation magnetization (Ms), coercivity (Hc), and retentivity (Mr). The values of Ms and Hc varied anisotropically with Sm doping. The findings of this study suggest that zinc cobalt ferrite doped with samarium might be beneficial for magnetic resonance imaging in biomedicine.

Keywords: Samarium, ferrite, nanoparticle; XRD, FESEM, HRTEM, FTIR, Electrical and magnetic properties

1. INTRODUCTION

Even many years after their development, ferrite materials still have great curiosity in the scientific community as well as commercial applications due to their remarkable properties in both nano and bulk form. Scientists, researchers, technologists, and engineers have put in a significant amount of work to discover new commercial uses for ferrites with extraordinary features [1, 2]. Because of its outstanding electrical resistivity, mechanical hardness, chemical

stability, and low cost, Co-Mg has piqued the curiosity of many researchers [3]. As a result of its exceptional characteristics, Co-Mg is an excellent candidate for use in a variety of applications, including microwave devices, memory devices, and transformer centers [4]. The mixed spinel structure $(\text{Mg}^{2+}\text{Fe}^{2+})[\text{Co}^{2+}\text{Fe}^{3+}]\text{O}_4$ was observed in Co-Mg ferrite [1], in which Mg^{2+} and Fe^{2+} occupy A site while Co^{2+} and Fe^{3+} occupy B site, has been discovered in this material. In addition to the technique and circumstances of preparation, as well as the quantity and kind of additives used, the characteristics of ferrites are impacted by many other factors. Due to their higher ionic radius, rare-earth implantation can fine-tune the structural, electrical, and magnetic properties of ferrites [5–11]. In ferrites, changing the cation dispersion between the A and B sites for divalent (M^{2+}) or ferric (Fe^{3+}) ions changes the cation dispersion. The structural, electrical, and soft magnetic properties of Co-Mg ferrites have been altered by the addition of substituted Ni, Si, and Mo atoms. [11-13]. Electron hopping between Fe^{2+} and Fe^{3+} is the main contributor to the electrical conductivity of spinel ferrites (generated through the sintering procedure [14, 15]) and Fe^{3+} ions at the octahedral (B) site [16]. The octahedral (B) site is occupied by rare-earth ions, which limits electron hopping by restricting the mobility of Fe^{2+} and increasing resistivity [8–10]. Some rare-earth ions, such as Gd^{3+} and Ho^{3+} , have previously been successfully replaced in the CoZn ferrites [9,10]. By substituting Sm^{3+} for Fe^{3+} in $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00 \leq x \leq 0.0025$ in a step of 0.005), the electrical resistivity of Co-Mg ferrites to compositions of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($x = 0.00 \leq x \leq 0.0025$ in a step of 0.005) has been expected to be increased, and it was expected that a significant characteristic (rise resistivity) would be obtained. Furthermore, in the ferrites Ni-Zn, Mg, Mg-Zn, and Co, the element samarium (Sm) was previously substituted [17–25]. To our knowledge, no research has been done on the effects of Sm replacement in Co-Mg ferrite. As a result, we present here the replacement of the Sm^{3+} ion for Fe^{3+} in Co-Mg ferrites with compositions of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ (where $x = 0.00 \leq x \leq 0.025$), which provides useful information for practical applications [26].

2. MATERIALS AND METHOD

Materials

The Sm-doped Co–Mg ferrite nanoparticles were made with analytical grade (AR) ingredients that did not require further purification. The ferrites were made using cobalt nitrate hexahydrate ($\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$), magnesium nitrate hexahydrate ($\text{Mg}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$), ferric nitrate hexahydrate ($\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$), and samarium nitrate hexahydrate ($\text{Sm}(\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$). The production of the Co–Mg ferrite nanoparticles was carried out using triple-distilled water.

Method

The samarium ion substitution these Co–Mg ferrite powders ($\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$) with a m/o of ($x = 0.000, 0.005, 0.010, 0.015, 0.020, 0.025$) were made using the citric sol–gel auto-combustion process. Citric acid was employed 1:3 with metal nitrate as a cation-exchange chemical. The stoichiometric weights of ($\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$), ($\text{Sm}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$), and ($\text{Mg}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$), were used to start the hydrolysis process ($\text{NO}_3)_3 \cdot 6\text{H}_2\text{O}$), ($\text{C}_6\text{H}_8\text{O}_7 \cdot \text{H}_2\text{O}$), and ($\text{Fe}(\text{NO}_3)_3 \cdot 9\text{H}_2\text{O}$), were all dissolved in distilled water and heated for 1 hour at 250°C with constant stirring. To keep the pH^7 , added the ammonia solution drop by drop. Using an electrically controlled magnetic stirrer and heater, the solution is allowed to heated to 150°C on a hot plate for 15 to 20 minutes. The solution is continuously stirred and heated, resulting in an extremely thick, viscous, dark-colored gel. Continue to heat the gel until all water

molecules have been removed, resulting in the start of the gel's combustion. Shortly after auto-ignition, the gel burns down to soft ash. The powder was grounded with an agate mortar for up to 3 hrs. In a temperature controlled furnace, the powder was annealed at 700°C up to 5 hours, until it was hard. The KBr pellet method was used to prepare pellets in the shape of discs with diameters of 15 mm and a thickness of 2 mm from ferrite powders. Silver paste was used on both sides of the pellets to achieve good ohm connections. These pellets were used in electrical and dielectric investigations.

Characterization

X-ray diffraction (XRD; Rigaku Miniflex-600, Japan) was used to characterise nanoparticles utilising CuK radiation ($\lambda = 1.5405$) from 20° to 80° and a scanning rate of 2 min. The surface morphology was examined using a FESEM, (FESEM; Oxford Instrument Inca X-act) and energy-dispersive X-ray (EDX) examination at a 200 kV accelerating voltage. A HRTEM was used to examine the cross-sections and particle sizes of the nanoparticles (HRTEM, JEM-F200 Multi-purpose Electron Microscope). The KBr pellet technique was used to measure Fourier-transform infrared (FTIR) spectroscopy using an IR spectrometer (FTIR-8400S Shimadzu). By tracing M-H curvatures, saturation magnetization, remanance, and coercive field for the powder sample were determined at room temperature using the vibrating sample magnetometer (VSM) model Lake Shore Cryotronics 8600 Series. The dielectric tests were completed with the help of the (Novocontrol GmBH, Alpha-A High Performance Frequency).

3. RESULTS AND DISCUSSION

3.1. XRD Analysis

The X-ray diffraction pattern is the same for all $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($0.000 \leq x \leq 0.025$) spinel ferrite nano particles made by the sol-gel auto combustion method. Sintered for 5 hours at 700°C, Co-Mg composite samples were created. The XRD method was used to determine the crystal structure and crystalline phase pattern of the substance. This is a highly useful approach for estimating the parameters of crystalline structures. XRD analysis revealed the most intense peak at $2\theta = 35$ for a cubic spinel structure. Identical to the peak of Co-Mg ferrite, all of the peaks with hkl are (111), (220), (311), (222), (400), (422), (511), (440), (440), and (533). The positions and miller indices of the peaks suggest that an FCC type spinel structure has formed. All of the peaks in the pattern are well defined and matches with the JCPDS file no 22-1086, indicating that the ferrites used in the preparation have an FCC spinel structure too [27, 28].

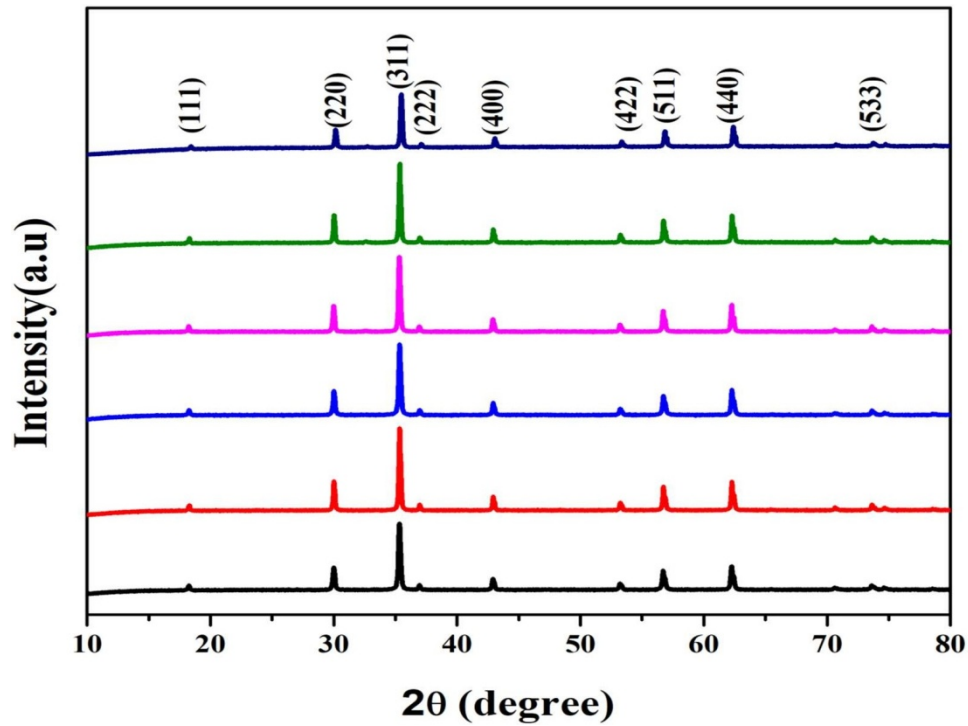


Fig.1. X-ray Diffraction pattern of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$).

Figure 1 shows the XRD patterns of the doped Co-Mg composition, and it is found that no impurity peaks were found. The XRD Debye Scherrer formula was used to compute the crystallite size (D_m) [28]:

$$D_m = \frac{k\lambda}{\beta \cos\theta} \text{ (nm)} \quad (1)$$

Where 'K' = 0.9 for spinel ferrites, ' λ ' = 1.5406 Å is the x-ray beam wavelength and ' β ' is the FWHM of the most intense peak and ' θ ' is the angle of the most intense peak. The crystallites size was estimated to be between 21.32 and 26.74 nm, which are absolutely small. Because of the increasing concentration of Sm^{3+} doping, the size of the crystallites rises, and they are now in this same range of 21.32 nm (0.000) to 26.74 nm (0.025) [29, 30]. The formula for calculating the average lattice constant is as follows:

$$a = d\sqrt{h^2 + k^2 + l^2} \text{ \AA} \quad (2)$$

The crystal plane spacing is 'd', and the miller index values are hkl. For all the samples, the average lattice constant ranges from 8.03Å to 8.89Å. The average lattice constant 'a' was computed using Sm^{3+} and Fe^{3+} ionic radii. It was observed that when the m/o of Sm^{3+} doping was increased, the average lattice constant increased as shown in Fig 1 [31].

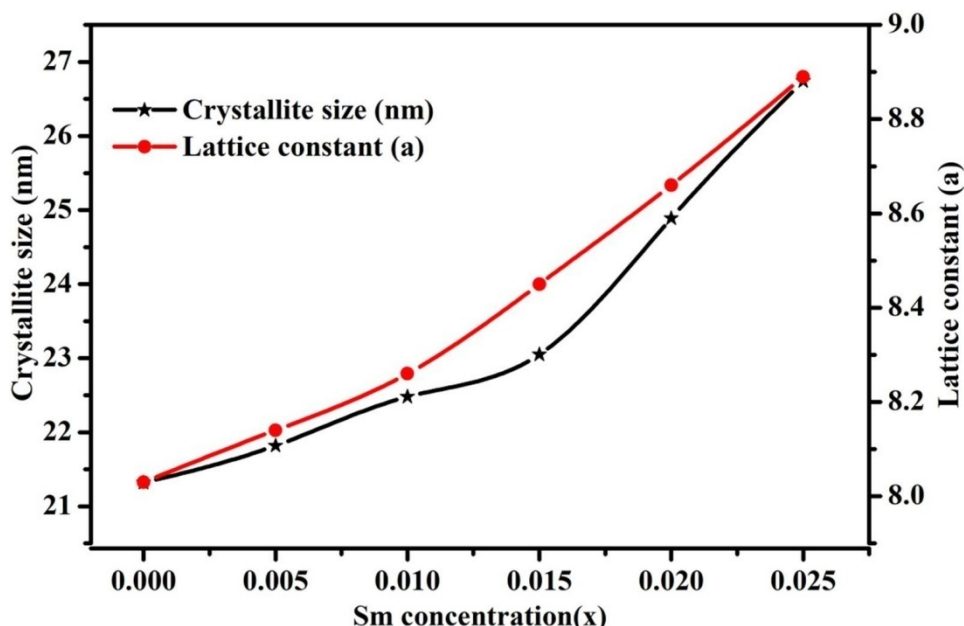


Fig.2.The variation of the Lattice constant and Crystallite size for $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$)

The formula for calculating the X-ray density ρ_x as follows:

$$\rho_x = \frac{8M}{Na^3} \text{ (g/cm}^3\text{)} \quad (3)$$

Where 'M' is the composition's molecular weight, 'Z' = 8 is the number of molecules per unit cell, 'N' is indeed the Avogadro number (6.0238×10^{23}), and for 'a³' is the volume of a unit cell. The density of X-rays was determined to be between 5.60 g/cm^3 to 5.99 g/cm^3 . The relationship between the density of X-rays and their concentration Sm^{3+} . As the quantity of Sm^{3+} doping rises, the X-ray density initially increases, then drops. These differences are due to the greater molecular weight of Sm^{3+} [32]. The bulk density of the pellets may be calculated using the relation:

$$\rho_m = \frac{m}{v} \quad (4)$$

The pellets' mass is 'm' and their volume is 'v'. On average, 3.43 g/cm^3 to 3.85 g/cm^3 were computed. The concentration of Sm^{3+} doping causes the bulk density to shows random change and a maximum of 0.025 m/o. As a result, lowering the bulk density lowers the overall weight of pallets [33].

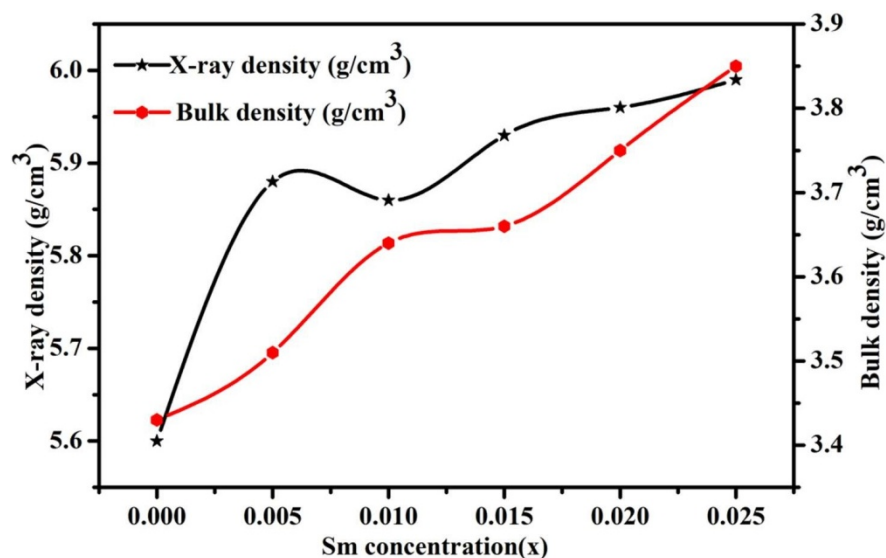


Fig.3. The Variation in the Bulk density and X-ray density of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$)

Table 1. Particle size (D_{Xrd}), Lattice parameter (a), Volume of unit cell (V), Strain (ϵ_s), jump length of A-site (L_A), jump length of B-site (L_B), $\text{Co}_{0.2}\text{Zn}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$).

Composition (x)	D_{Xrd} (nm)	a (Å)	V (Å ³)	L_A (Å)	L_B (Å)	Strain (10^3)	Dislocation density
0.000	21.32	8.03	598.14	3.54	2.79	25.20	9.823
0.005	21.82	8.14	598.26	3.66	2.91	22.61	8.481
0.010	22.48	8.26	598.31	3.36	2.69	20.10	7.298
0.015	23.05	8.45	598.49	3.86	2.79	18.24	6.250
0.020	24.89	8.86	598.62	3.26	2.59	17.18	6.433
0.025	26.74	8.89	598.81	3.816	2.88	14.20	5.340

Table 2. Bulk density (ρ_B), X-ray density (ρ_x), porosity (P), Absorption band wave number (V_1 and V_2) From FTIR spectra, for nanosized $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$).

Composition (x)	ρ_x (g/cm ³)	ρ_B (g/cm ³)	P (%)	V_1 (cm ⁻¹)	V_2 (cm ⁻¹)
0.000	5.60	3.43	14.001	556	410
0.005	5.88	3.83	9.134	548	401
0.010	5.86	3.54	9.512	553	395
0.015	5.93	3.66	8.587	559	409
0.020	5.96	3.75	8.901	552	403
0.025	5.99	3.85	7.583	555	427

Using d_x and d_B values, porosity (P) values were calculated and found to be in the range between 14 to 7 percent as the doping m/o is increased (Table 2), indicating that the synthesised materials are porous [34-36].

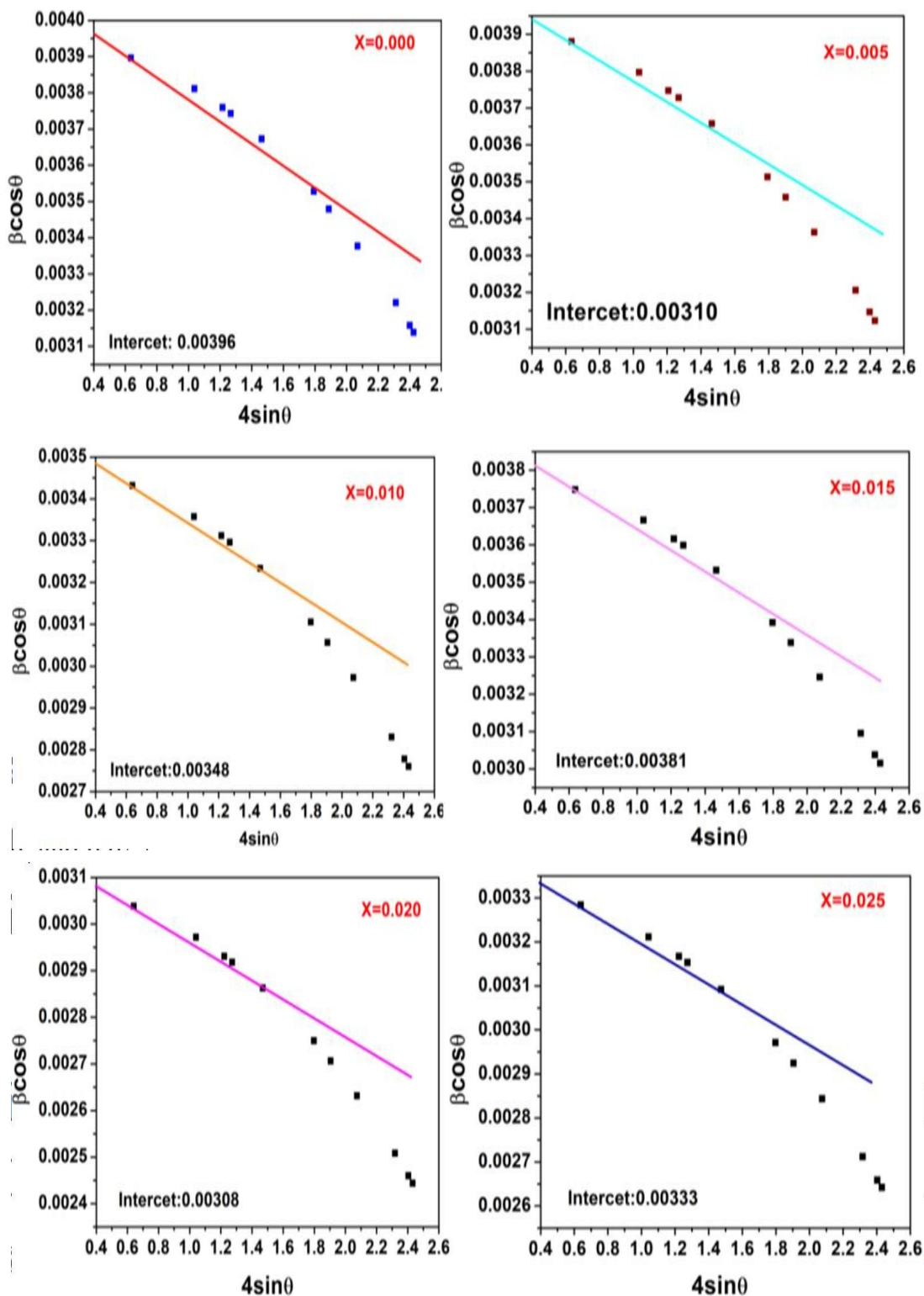


Fig.4. Williamson-Hall plot of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites $x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025 .

The following formula is used to compute the strain:

$$\varepsilon = \frac{\beta}{4 \tan \theta} (10^{-3}) \quad (5)$$

Where ' β ' is the diffraction angle of the most intense peak. The predicted strain was 14.01×10^{-3} to 7.58×10^{-3} psi. As the quantity of Sm^{3+} doping is increased, the strain reduces randomly. At a value of $x = 0.000$, the highest amount of strain was found [37]. To calculate the dislocation density of the nanoparticles produced.

$$\delta = \frac{1}{D^2} \quad (6)$$

Where 'D' is the crystalline size. With increasing m/o of Sm^{3+} concentration, the dislocation density decreases at minimum for $x = 0.000$ to 0.025 is 9.82 to 5.34 line/m. The jump length L_A and L_B between both the ions at A and B-site were calculated using the relationship.

$$L_A = \frac{a\sqrt{3}}{4} \quad (7)$$

$$L_B = \frac{a\sqrt{2}}{4} \quad (8)$$

In Co-Mg ferrite, the values of jump length L_A are increasing and L_B are decreased with Sm^{3+} increasing composition. This is owing to the fact that the lattice constant lowers with the addition of Sm^{3+} composition [38-40].

4. FIELD EMISSION SCANNING ELECTRON MICROSCOPE (FESEM)

Figure 5 depicts the surface morphology or structural investigation of ferrite samples at room temperature using a FESEM for ruptured surfaces. Sm doped $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ferrites FESEM are shown in Fig 5. According to the high magnification picture at 200 nm, this shows agglomerated nanoparticle morphology, which is composed of secondary irregularly shaped spherical nanoparticles. These nanoparticles are 70 nm–120 nm in size [41]. The presence of Sm-doping in $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ ($0.000 \leq x \leq 0.025$) results in increased agglomeration of the spherically nanoparticles, which was accompanied by an increase in the crystalline size [42]. Higher concentrations of Sm agglomerate owing to ferromagnetic Co^{2+} and diamagnetic Mg^{2+} replacement. The magnetic character of the materials is to blame for the agglomeration, which was further verified by VSM analysis [43].

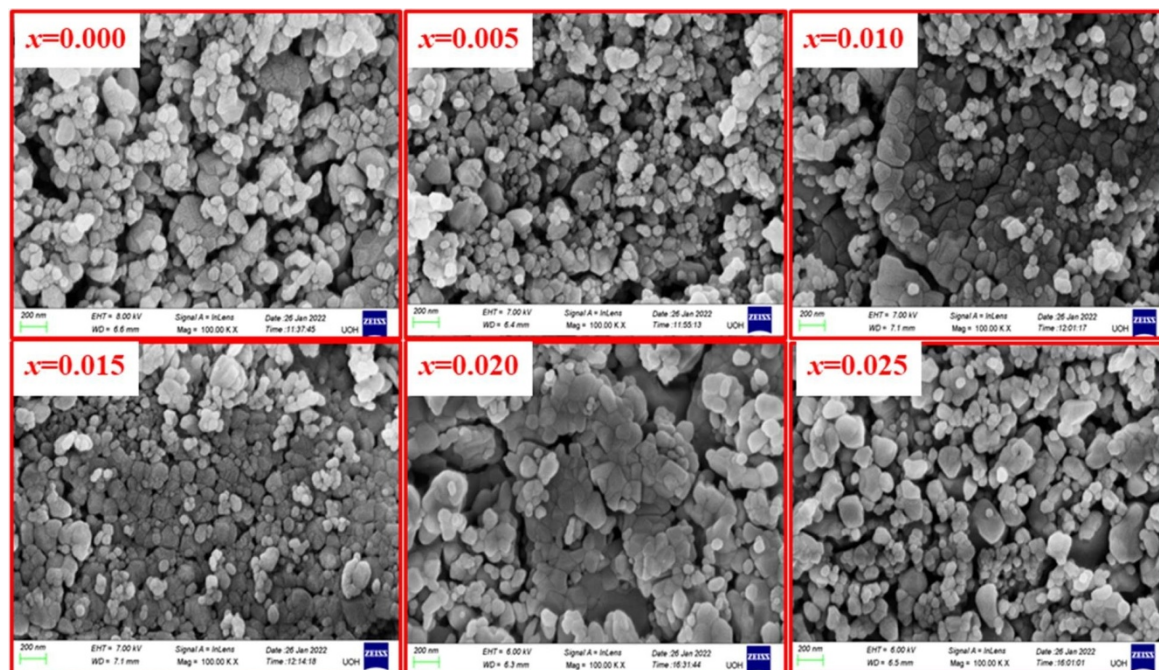


Fig. 5. FE-SEM Images of the $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

5. HIGH-RESOLUTION TRANSMISSION ELECTRON MICROSCOPY (HRTEM)

TEM pictures of the $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ sample are shown in Fig 6. The majorities of nanoparticles have a spherical shape and are seen agglomerated. It is possible that single-phase crystals were formed, as shown by the presence of strong diffraction rings in TEM micrographs [44]. ImageJ software calculated the average particle size to be 74nm. Figure 6 shows the (SAED) pattern of the nanoferrite CMSFO ($x = 0.000, 0.015, 0.025$), which displays patchy diffraction rings that exactly match the diffraction peak obtained from the XRD pattern (see Fig.7), confirming the cubic spinel phase development in Sm^{3+} doped Co-Mg nanoferrites. Fig. 6 shows the HRTEM of the nanoferrite ($x = 0.000, 0.015, 0.025$). Several lattice fringes can be noticed. Fringe widths of 0.245, 0.264, and 0.269 nm correspond to the (311) peak, demonstrating that HRTEM and XRD investigation of Co-Mg-Sm nanoferrites are identical.

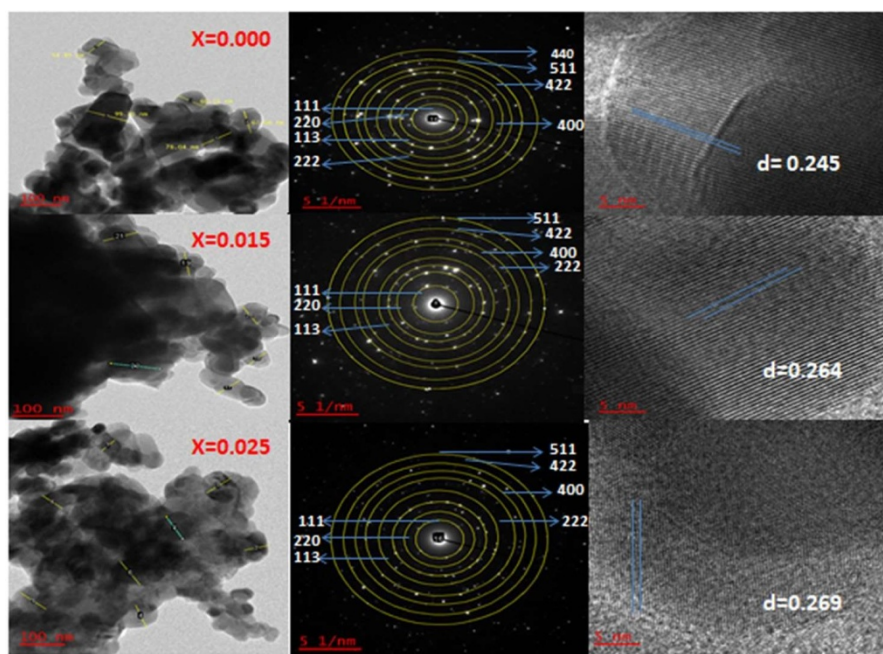


Fig. 6. Show the TEM Micrographs of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$; (where $0.000 \leq x \leq 0.025$, step 0.005) SAED and HRTEM of nanoferrite with $x=0.000, 0.015, 0.025$

6. FOURIER TRANSMISSION INFRARED SPECTROSCOPY (FTIR)

The FTIR spectra of nanosized $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ in the $350\text{--}4000\text{cm}^{-1}$ range were acquired at room temperature and are shown in Figure 7. The FTIR spectra of all materials exhibit two primary absorption bands in the $350\text{--}700\text{cm}^{-1}$ region. The inherent stretching vibrations of the metal at the tetrahedral $\text{M}_{\text{tet}}\text{--O}$ and octahedral $\text{M}_{\text{oct}}\text{--O}$ sites are related to these two vibration bands. Tetrahedral complexes are represented by the band that appears at the higher wave number (ν_1 , $548\text{--}560\text{cm}^{-1}$), whereas octahedral complexes are represented by the band that appears at the lower wave number (ν_2 , $395\text{--}430\text{cm}^{-1}$) [45]. By adding Sm^{3+} content, it is possible to determine that the difference in the $\text{Fe}^{3+}\text{--O}^2$ distance for the octahedral and tetrahedral sites is what causes the difference between the ν_1 and ν_2 band locations [46]. In $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ doped with the rare earth element, the ν_2 band broadens, indicating the occupancy of rare earth ions on the B-sites [47]. Stretching vibrations of the C=O are shown by the peaks between 1630 and 1660cm^{-1} . A stretched hydroxyl (O-H) group that corresponds to a bandwidth between 3435 and 3480cm^{-1} is seen, signifying the presence of moisture and water. The observed peak between 2330 and 2370cm^{-1} is caused by the stretching of the C-H bond [48].

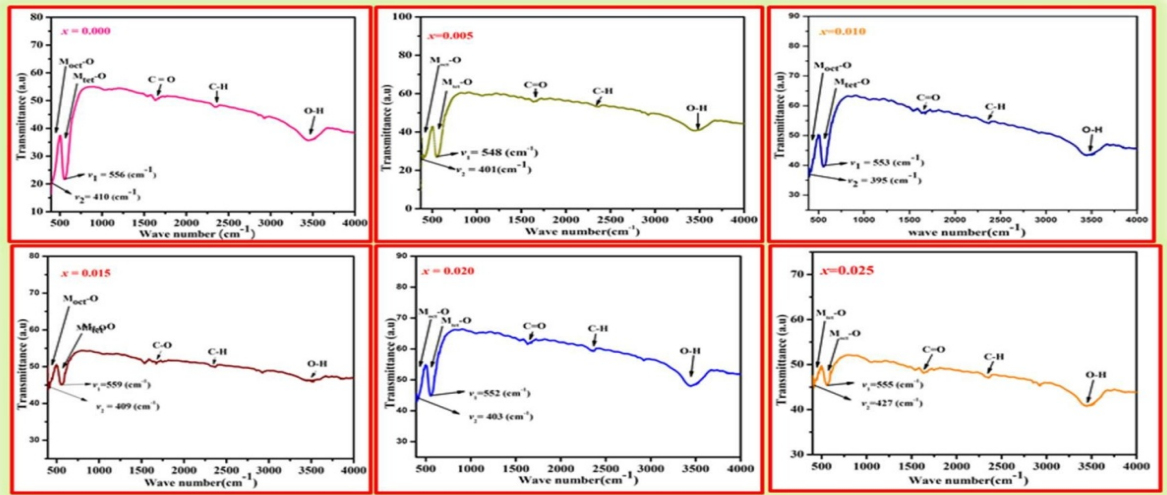


Fig.7. FTIR spectra of nanosized $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

7. ELECTRICAL PROPERTIES

7.1 Dielectric Measurement

The dielectric dispersion of the samples was examined from 3Hz to 20MHz at room temperature. As shown in Figure 8, a frequency applied field produces Sm-doped Co-Mg ferrite nanoparticles. Interfacial polarisation response of dielectric dispersion decreased with applied field frequency [49-52]. It decreased rapidly up to 7 KHz and reaches to saturation. Because the domains in a dielectric material were initially random, adding a field oriented them in the field's direction.

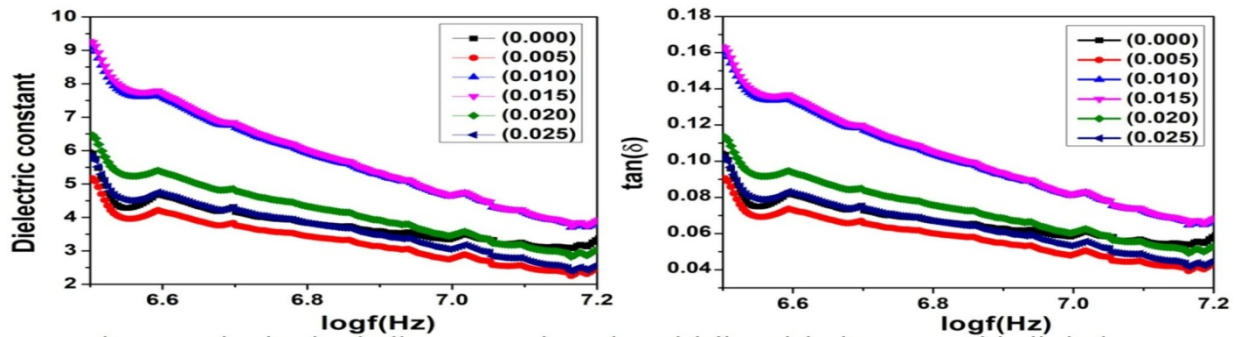


Fig.8. dielectric constant spectra of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

Initially, the domains aligned with the applied field and responded well to switching. In the study, the domains lagged behind the applied field by over 7MHz. Increasing the amount of Sm doping also increased the dielectric constant, proving that the hopping of charges between charge carriers, originally between Fe^{3+} and Fe^{2+} , was successful [53].

Fig 8 represent the frequency dependent dielectric constant for CMSFO systems in the range of 3Hz to 20Hz from 700oC. It is observed that dielectric constant is high at lower frequencies and decreases exponentially reaching to constant value at very high frequencies for all Sm doped systems. This is due to the presence of space charge polarization along with other polarizations such as dipolar, ionic and electric. Larger value of dielectric constant for doped system can also be due to the increased number of lattice defects compared to that of pure system. The increase in dielectric with m/o is due to interfacial polarization, grain boundary defect.

7.2 Dielectric loss ($\tan\delta$)

Dielectric loss is a critical metric in determining a material's viability for electrical applications, notably in capacitors and bridge circuits. Any electronic application requires a dielectric loss of less than unity, suggesting suitability for high resistance applications and ease of manufacture [54]. Figure 10 shows dielectric loss in Co–Mg nanoparticles with Sm doping from 3Hz to 20 MHz. The dielectric loss had a typical response to the applied field frequency, with the loss reducing as the frequency was increased. The dielectric loss is maximum for 0.164

The variation of dielectric loss (tan) with frequency at various temperatures of CMSFO system is shown in Fig 8. It can be noticed that the dielectric loss is maximum at the lower frequencies and found to decrease gradually with frequency. The dielectric loss represents the energy loss that occurs when the polarization lags behind the applied electric field due to the presence of grain boundaries [55]. The loss in these systems could be attributed to various conduction processes. The decrease in dielectric loss with frequency can be explained on the basis of Koops theory. It can also be noticed from Fig that change in dielectric loss with m/o sm and was observed maximum for x=0.015 system [56].

7.3 AC Conductivity:

The frequency dependent AC conductivity for different systems was studied at 700C shown in Fig 9. Study of AC conductivity give the electrical transport properties of a material by using jonschers power law [57];

$$\sigma_{AC}(\omega) = \sigma_{DC} + A\omega^n \quad (9)$$

Where σ_{DC} is frequency independent conductivity at low frequencies. Random diffusion of the charge via activated hopping gives rise to a frequency independent conductivity. At higher frequencies, σ_{AC} exhibits increases in dispersion and eventually becoming almost linear. A is the temperature dependent pre-exponential factor and n is the frequency independent exponent, but depends on temperature and material. The increasing value of σ_{AC} with frequency is due to the hopping of cations between the neighboring of sites. Also, increases of conductivity with temperature could be due to the thermally activated processes. From the Fig 9 the enhancement of conductivity was observed maximum for 9.027 m/o, high could be due to presence of large number of charge carriers and the increases rate of hopping, there after a decreases in the conductivity was noticed for further increases in Sm concentration which could be due to decreases in grain density or the formation of neutral complexes. For 5.700m/o, conductivity is lower than that pure of pure system [58-60].

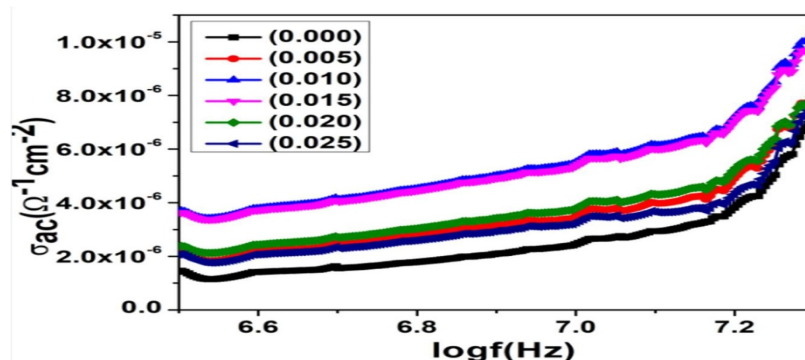


Fig.9. AC conductivity spectrum of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020, \text{ and } 0.025$).

7.4. Real Impedance and Imaginary Impedance

Frequency dependent real and imaginary impedance spectra of frequency from 3Hz to 20MHz is shown in Figure 10 and 10 (b). See also real (Z') and imaginary impedance (Z'').

$$Z' = |Z| \cos\theta z \quad (10)$$

$$Z'' = |Z| \sin\theta z \quad (11)$$

In all compositions, inter-facial polarization causes resistive grain barriers, resulting in high resistance. Less real and imaginary impedances mean more blending of impedance curves at higher frequencies [61, 64].

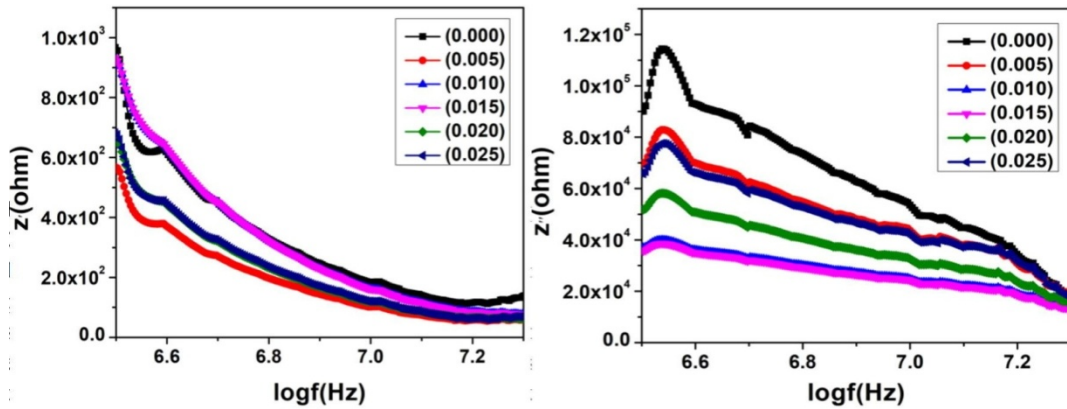


Fig.10. Variation of real part and imaginary part of impedance with frequency spectrum of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

Impedance spectroscopy analysis was carried out to study the electrical and dielectrical properties of the material in the frequencies between 3Hz and 20MHz. It can also be used to separate the contributions of electrical properties due to the grains, grain boundaries and electrode/ electrolyte interface in a polycrystalline material. Fig shows the impedance spectra of CMSFO system with varying m/o at 700°C . Frequency dependent real and imaginary impedance response shown in Fig. For real impedance graph shows that at low frequency ranges the impedance is maximum, thereafter it appears to merge in the high frequency region which could be due to the relaxation of space charge, dipolar polarizations. For imaginary impedance the Sm doped Co-Mg-Fe is maximum for 9.0×10^4 , $x = 0.000$, Imaginary impedance is also similar to real and is the value is maximum for 9.5×10^2 , $x = 0.015$ [65].

8. VIBRATING SAMPLE MAGNETOMETER (VSM)

Figure 11 shows nanosized $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ M-H loops ($x = 0, 0.00 \leq x \leq 0.025$). In M-H loops at ambient temperature, saturation magnetization (M_s), coercivity (H_c), retentivity (M_r), magnetic moment (B), squareness ratio (SQR), and anisotropy constant (K) have all been calculated. The obtained values are presented in Table 3 for the range of -18 to 18 Oe. The resulting saturation (M_s) in the range between 96.58 - 46.40 emu/g, M_r , H_c , and SQR, values are lowered, whereas H_c rises as the amount of Sm^{+3} in the mixture increases. The loops for the produced samples evaluated at room temperature are relatively near to one another, suggesting extremely strong coercivity and typical super paramagnetic activity, as shown in Figure 11. The decrease in M_s with increase in particle size might be explained by cation

redistribution (exchange and substitution of cations in octahedral and tetrahedral sites) [66, 70]. The structure of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ transforms from an inverse spinel to a mixture spinel as the cations disperse. M_s decreases when $x=0.005$ until the nanoferrite with after which M_s slightly rises for the nanoferrite with $x=0.010$, and then M_s decreases again when $x=0.025$ is employed. The existence of a magnetically dead layer on the surface of the material, as well as a degree of spin canting over the whole volume of the crystallite structure, might explain the observed decrease in M_s . The quick rise in M_s for the nanoferrite with $x=0.010$ is due to an increase in the size of its crystallites, while the rapid increase in M_s for the nanoferrite with $x=0.010$ is due to an increase in the size of its crystallites [71-73]. A-A, B-B, and A-B super-exchange interactions, as well as A-B super-exchange interactions, are seen in spinel ferrites. The cation distribution between A and B varies as Sm^{3+} dopes the ferrite, as does the intensity of the A-B exchange contact. A- sites get less Fe^{3+} than B- sites when Sm^{3+} ions are introduced. Sm^{3+} most likely enters the octahedral site because to its large radius and tight tetrahedral A-site gap. The reduction in Fe^{3+} concentration reduces the B-B exchange contact at the octahedral site, generating canting spins at the nanoparticles surface. According to table 3, the saturation magnetization decreases when the net magnetic moment $B=M_B-M_A$ decreases (M_A and M_B are the Band A sublattice magnetic moments). When the concentration of Sm^{3+} in CMSFO compositions is increased, the crystal lattice anisotropy decreases, resulting in a reduction in M_r [74]. Anisotropy increases until $x=0.0010$, after which it decreases constantly until $x=0.025$. The value of H_c increases as the concentration of Sm^{3+} increases. This is owing to the permutation of Co^{2+} ions with high magneto-crystalline anisotropy but low Sm^{3+} ion concentration. In fact, in addition to having non-quenched orbital angular momentum, Co^{2+} ion also possess a spin-orbit coupling. When spin-orbit coupling grows, magneto crystalline anisotropy rises, and materials containing cobalt ions become more coercive [75]. The corecivity of the nanoferrite with $x=0.010$ develops rapidly due to its high M_s Value. To overcome this spin alignment, a high value of H_c must be used. Because of the varied concentrations of Sm^{3+} , the squareness ratio decreases suddenly; it decreases for $x=0.005$, increases for $X=0.010$, and then decreases again for $X=0.025$ [76.77].

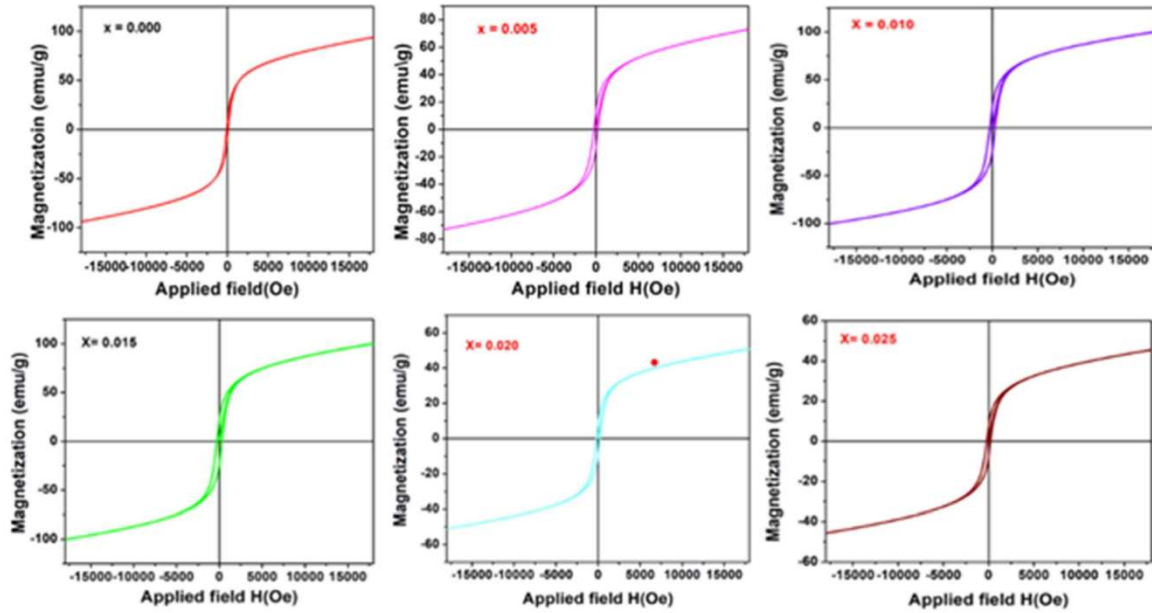


Fig11. M-H curves of nanosized $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

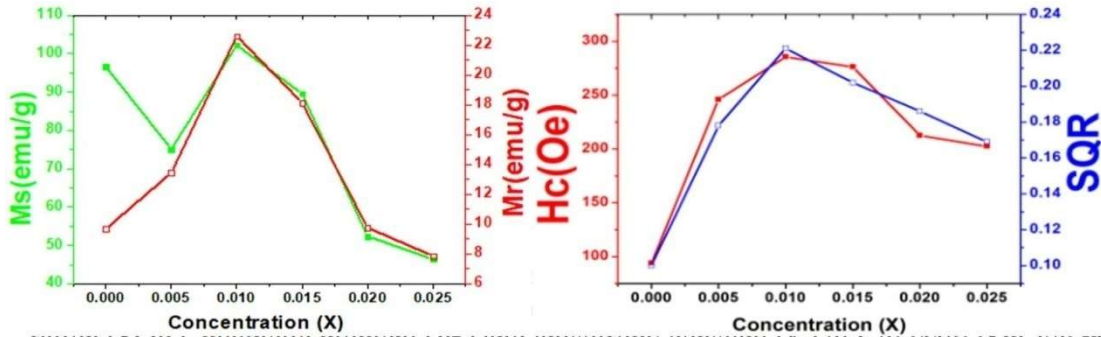


Fig.12. Saturation magnetization (M_s), retentively (M_r), and squareness ratio H_c [Oe], coercively (SQR) spectrum of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$ nano-ferrites ($x = 0.000, 0.005, 0.010, 0.015, 0.020,$ and 0.025).

Table3. The saturation magnetization (M_s), retentively (M_r), coercivity (H_c), squareness ratio (SQR), magnetic moment (η_B) and anisotropy constant (K) of $\text{Co}_{0.2}\text{Mg}_{0.8}\text{Sm}_x\text{Fe}_{2-x}\text{O}_4$; (Where $0.000 \leq x \leq 0.025$, step 0.005) nanoferrites.

Sample code	$M_s(\text{emu/g})$	$M_r(\text{emu/g})$	$H_c(\text{Oe})$	SQR	$\eta_B(\mu_B)$	$K(\text{erg/cm}^3)$
0.000	96.58	9.67	94.12	0.100	4.146	9275.622
0.005	75.11	13.44	246.08	0.178	3.321	18860.274
0.010	102.24	22.61	285.57	0.221	4.406	29792.527
0.015	89.54	18.14	276.26	0.202	3.866	25241.143
0.020	52.41	9.77	212.52	0.186	2.267	11365.482
0.025	46.40	7.86	202.46	0.169	2.011	9585.861

9. CONCLUSION

In this study, the citric sol gel auto combustion process is used to make samarium (Sm) substituted cobalt zinc ferrite ($\text{CoMgFe}_2\text{O}_4$) (range) nanoparticles. Because it is a straight forward and direct procedure. It can also be done in a natural setting. The mixing of solutions, stirring, drying, grinding, and annealing are the essential phases in this process. At 700°C , all of the prepared samples were sintered. Many studies have looked at the cations distribution of Sm- substituted ($\text{CoMgFe}_2\text{O}_4$) ferrite, and it has been found to be temperature dependent. The samples were found to be nicely crystalline spinel ferrite, which was regarded as a single phase ferrite, according to the XRD calculations. The crystallite size ranged from 21.32 to 26.74 nm, and the average lattice parameter was 8.03-8.89 (according to the XRD parameters). This is a highly useful approach for estimating crystalline characteristics, including strain, hopping length (L_A - L_B) site, dislocation density, bulk density, and X-ray density porosity. The XRDs most intense peak was reported at $2\theta=35^\circ$. The peak's locations and miller indices indicate the formation of an FCC spinel structure. According to the FESEM results, the particles are evenly dispersed, homogenous, and somewhat agglomerated. A TEM investigation revealed perfectly spherical and homogeneous nanocrystalline particles. TEM studies confirmed the nanoparticle sizes (74-80nm). A powerful ring of light surrounds the spinel ferrite sample in the SAED pattern. FTIR studies indicated two substantial peaks spanning from 395 cm^{-1} to 560 cm^{-1} , showing metal-oxygen stretching and the formation of ferrite composites. The material has Maxwell- Wagner interfacial polarization, which decreases as the frequency of the applied field rises, according to dielectric spectroscopy. With a dielectric loss of less than 8, the sample might be utilised to make multilayer inductors. The g-value was decreased and the spin-spin relaxation and the resonance magnetic field were increased by increasing the concentration of Sm^{3+} . The samples' Ac conductivity rose as doping and frequency were increased. The frequency window related to complex impedance appeared for $x = 0.015$ and $x = 0.000$ at room temperature in the frequency. At ambient temperature, saturation magnetization, coercivity, and remanent magnetization from M-H loops were lowered when Sm^{3+} was doped in zinc cobalt ferrite. The M-H loop analyses show the super paramagnetic behaviour of the prepared samples.

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“DECLARATION OF COMPETING INTERESTS”

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

“RESEARCH DATE POLICY AND DATA AVAILABILITY STATEMENT”

Data available at the corresponding author will be provided on the request.

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